

Math 5440
Wed 11/24

Eigenfunctions of the Laplacian (and other self adjoint operators) (connects to chapters 6-7 in Weinberger)

Overview:

Recall for wave & heat eqns in arbitrary space dimension

- Duhamel \Rightarrow solve homog PDE yields integral formulas for non-homog. PDE
- If you try separation of variables in bounded domains you end up wanting to use eigenfunctions of Δ (or $-\Delta$), with appropriate Dirichlet or Neumann BC's:

$$u_t - k \Delta u = 0$$

if $u = X(x)T(t)$

$$T'X - kT\Delta X = 0$$

$$\frac{\Delta X}{X} = \frac{T'}{kT} = -\lambda$$

$$u_{tt} - c^2 \Delta u = 0$$

if $u = X(x)T(t)$

$$T''X - c^2T\Delta X = 0$$

$$\frac{\Delta X}{X} = \frac{T''}{c^2T} = -\lambda$$

notice $\lambda \geq 0$ so $-\lambda \leq 0$ because if $X(x)$ solves $\Delta X = -\lambda X$

then $\int_{\Omega} |\nabla X|^2 + \underbrace{X \Delta X}_{-\lambda X^2} = \int_{\partial \Omega} X \frac{\partial X}{\partial n}$

so $\lambda = \frac{\int_{\Omega} |\nabla X|^2}{\int_{\Omega} X^2}$ " 0 for NP & DP

thus your separated solutions will be have as you expect (exponentially decaying or const for HE, waves for wave eqn.)

- to solve IBVP's with DP or NP

HE $u(x,0) = f(x) \quad x \in \Omega$

WE $u(x,0) = f(x)$
 $u_t(x,0) = g(x)$

need to know whether the eigenfunctions are "complete", and maybe other stronger convergence properties.

Theorem Let $\Omega \subset \mathbb{R}^n$, bounded, piecewise smooth boundary.

Then there is an orthonormal complete collection of Dirichlet eigenfunctions

$$\{\psi_k\}_{k \in \mathbb{N}}$$

$$-\Delta \psi_k = \lambda_k \psi_k$$

$$0 < \lambda_1 < \lambda_2 \leq -\lambda_3 \leq \lambda_4 \rightarrow \infty$$

(Similarly for Neumann eigenfunctions)
 $0 = \mu_1 < \mu_2 \leq \mu_3 \dots$

with respect to $\langle f, g \rangle = \int_{\Omega} fg \, dV_n$

$$\|f - \text{proj}_{V_N} f\| \rightarrow 0 \text{ as } N \rightarrow \infty.$$

We'll sketch this proof, as a special case of the spectral theorem for compact self adjoint operators on Hilbert Spaces

$$L: V \rightarrow V$$

$$\langle u, Lv \rangle = \langle Lu, v \rangle$$

$$\forall u, v \in X.$$

complete inner product spaces

separable

\exists countable dense subset.

L is compact (linear)

transformation, $L: V \rightarrow V$, if $\forall \{f_k\} \in V$, a bounded sequence ($\|f_k\| \leq M \forall k$) the image sequence $\{L(f_k)\}$ has a convergent subsequence.

(If V is finite dim'd and L is linear it's automatically compact.)

Remark: Our $V = L^2(\Omega) = \{f: \Omega \rightarrow \mathbb{R} \text{ s.t. } \int_{\Omega} f(x)^2 < \infty\}$.

Our $L \neq \Delta$, which isn't even defined on all of V .

Rather it's " Δ^{-1} ", i.e. for $F \in L^2(\Omega)$ $\Delta_{DP}^{-1}(F)$ is

$$\text{the soltn to } \begin{cases} \Delta u = F & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases}$$

as given either by the variational or Green's function formulations.

Examples we know of self adjoint operators.

- \mathbb{R}^n , $L(\vec{x}) = A\vec{x}$, A symmetric, usual dot product
 $\langle Ax, y \rangle = (Ax)^T y = x^T A^T y = x^T A y = \langle x, Ay \rangle$.

- Δ is self-adjoint on the subspace of $L^2(\Omega)$ given by
 $V_{DP} = \{u \in L^2(\Omega) \text{ s.t. } u \in C^2(\Omega) \cap C^1(\bar{\Omega}) \text{ and } u=0 \text{ on } \partial\Omega\}$
 $V_{NP} = \{u \in L^2(\Omega) \text{ s.t. } u \in C^2(\Omega) \cap C^1(\bar{\Omega}) \text{ and } \frac{\partial u}{\partial n} = 0 \text{ on } \partial\Omega\}$

because

$$\langle u, Lv \rangle - \langle Lu, v \rangle = \int_{\Omega} u \Delta v - \Delta u v = \int_{\partial\Omega} u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} = 0 \text{ if } u, v \in V_{DP} \text{ or if } u, v \in V_{NP}.$$

examples of complete eigenfun collections

- $\{\sin nx\}_{n \in \mathbb{N}}$ DP eigenfun of Δ on $[0, \pi]$
 $\langle fg \rangle = \frac{2}{\pi} \int_0^{\pi} f(x)g(x) dx$
- $\{\cos nx\}_{n \in \mathbb{N}} \cup \{\frac{1}{\sqrt{2}}\}$ NP eigenfun for Δ .

- $L(u) = \frac{1}{\rho(x)} \left[\frac{d}{dx} \left(p(x) \frac{du}{dx} \right) + q(x)u \right]$ for funs defined on the interval $\alpha \leq x \leq \beta$
 $\rho > 0$ and $\langle u, v \rangle = \int_{\alpha}^{\beta} u(x)v(x)\rho(x) dx$.

(these sorts of L arise when we consider disk, ball or cylinder domains and separate variables for Δ .)

$$\langle Lu, v \rangle = \int_{\alpha}^{\beta} \frac{1}{\rho} [(pu') + qu] v \rho = [pu'v]_{\alpha}^{\beta} + \int_{\alpha}^{\beta} quv - pu'v' dx$$

$$\langle u, Lv \rangle = \int_{\alpha}^{\beta} \frac{1}{\rho} [(pv') + qv] u \rho = [pv'u]_{\alpha}^{\beta} + \int_{\alpha}^{\beta} quv - pv'u' dx$$

④

Spectral Theorem: Let $(V, \langle \cdot, \cdot \rangle)$ be a Hilbert space

$L: V \rightarrow V$ linear and compact, and self-adjoint.

Then \exists a complete orthonormal basis $\{u_k\}_{k \in \mathbb{N}}$ of eigenvectors of L .

$$L u_k = \lambda_k u_k.$$

$$\lambda_k \rightarrow 0 \text{ as } k \rightarrow \infty.$$

proof on Monday.

HW for Fri Dec. 3

6.31.3 (double sine series in a square)

6.32.1, 4, 5 (solns to homog. Laplace, heat, wave eqns in rectangles)

Ⓘ Also, find the fully separated solutions to heat and wave equation in the unit disk, i.e.

$$u(\vec{x}, t) = R(r) \Theta(\theta) T(t)$$

$$\text{for } u_t - k \Delta u = 0$$

(relates to §7.41)

$$u_{tt} - c^2 \Delta u = 0$$

Do not try to solve the ODE's for $R(r)$

Ⓙ Same question in the unit ball, using spherical coords

$$x = r \sin \theta \cos \phi$$

$$y = r \sin \theta \sin \phi \quad \leftarrow \text{not like in m.v. calc.}$$

$$z = r \cos \theta$$

Step 1: derive $L = \Delta$ in spherical coords.

