

① Consider polar coordinates in  $\mathbb{R}^2$  for  $r \in \mathbb{R}$  with  $r = \sqrt{x^2 + y^2}$

$$\theta = \begin{cases} \arctan(y/x) & \text{I, IV} \\ \arccos \frac{x}{\sqrt{x^2+y^2}} & \text{I, II} \\ \text{etc.} \end{cases}$$

a) Show  $\ln r = \ln \sqrt{x^2 + y^2}$  is harmonic.

Proof:  $\ln r = \ln(x^2 + y^2)^{1/2} = \frac{1}{2} \ln(x^2 + y^2)$

$$\frac{\partial}{\partial x} \left[ \frac{1}{2} \ln(x^2 + y^2) \right] = \frac{2x}{2(x^2 + y^2)} = \frac{x}{x^2 + y^2}$$

$$\frac{\partial}{\partial y} \left[ \frac{1}{2} \ln(x^2 + y^2) \right] = \frac{2y}{2(x^2 + y^2)} = \frac{y}{x^2 + y^2}$$

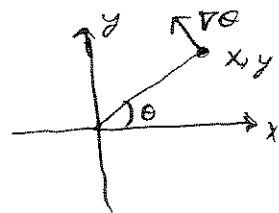
$$\frac{\partial^2}{\partial x^2} \left[ \frac{1}{2} \ln(x^2 + y^2) \right] = \frac{(x^2 + y^2) - x(2x)}{(x^2 + y^2)^2} = \frac{-x^2 + y^2}{(x^2 + y^2)^2}$$

$$\frac{\partial^2}{\partial y^2} \left[ \frac{1}{2} \ln(x^2 + y^2) \right] = \frac{(x^2 + y^2) - y(2y)}{(x^2 + y^2)^2} = \frac{x^2 - y^2}{(x^2 + y^2)^2}$$

$$\text{so } \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) \left[ \frac{1}{2} \ln(x^2 + y^2) \right] = \frac{-x^2 + y^2 + x^2 - y^2}{x^2 + y^2} = 0. \quad \square$$

b) Derive the formula below for  $\nabla \theta$  without using inverse trig functions:

$$\nabla \theta = \frac{1}{x^2 + y^2} \langle -y, x \rangle,$$



Soln: We know that  $\nabla \theta$  points in the direction of largest increase of  $\theta$ . This is in the direction shown in the figure, i.e. perpendicular to the radial direction (the direction in which  $\theta$  does not increase at all). Since  $\theta$  increases as we move counterclockwise,  $\nabla \theta$  will be in the direction of the counterclockwise rotation of  $\langle x, y \rangle$ , i.e.  $\nabla \theta$  will be in the direction of the unit vector

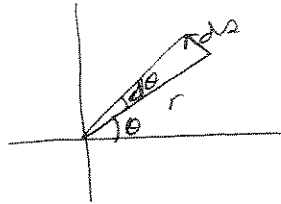
$$\hat{u} = \frac{\langle -y, x \rangle}{\sqrt{x^2 + y^2}}$$

$$\text{Thus, } \nabla\theta \cdot \hat{u} = \|\nabla\theta\| \|\hat{u}\| \cos(0) = \|\nabla\theta\|,$$

since  $\nabla\theta$  points in the same direction as  $\hat{u}$ .

$$\text{Now, } \nabla\theta = \|\nabla\theta\| \hat{u} = (\nabla\theta \cdot \hat{u}) \hat{u}.$$

(consider the figure below, where  $d\theta$  is an infinitesimal change in  $\theta$ ,  $ds$  is an infinitesimal change in arclength,  $\rho$  (approximated by the straight line  $ds$  drawn).



$$ds = r d\theta \Rightarrow \frac{d\theta}{ds} = \frac{1}{r}. \quad \text{But } \hat{u} \text{ is in the same direction as } ds,$$

$$\text{so } \|\nabla\theta\| = \nabla\theta \cdot \hat{u} = \frac{d\theta}{ds} = \frac{1}{r}. \quad \text{Thus, since } \nabla\theta = \|\nabla\theta\| \hat{u}$$

$$= \frac{1}{r} \hat{u} = \frac{1}{\sqrt{x^2+y^2}} \frac{\langle -y, x \rangle}{\sqrt{x^2+y^2}} = \frac{1}{x^2+y^2} \langle -y, x \rangle. \quad \square$$

c) Using the formula for  $\nabla\theta$  in b, show  $\theta$  is harmonic.

$$\text{Proof: } \Delta\theta = \nabla \cdot (\nabla\theta) = \frac{\partial}{\partial x} \left( \frac{-y}{x^2+y^2} \right) + \frac{\partial}{\partial y} \left( \frac{x}{x^2+y^2} \right)$$

$$= \frac{(x^2+y^2)(0) - (-y)(2x)}{(x^2+y^2)^2} + \frac{(x^2+y^2)(0) - x(2y)}{(x^2+y^2)^2}$$

$$= \frac{2xy}{(x^2+y^2)^2} - \frac{2xy}{(x^2+y^2)^2} = 0. \quad \square$$

d) Show that  $\Delta u = u_{xx} + u_{yy} = \frac{1}{r} \frac{\partial}{\partial r} (r u_r) + \frac{1}{r^2} u_{\theta\theta}$ .

By the chain rule,

$$u_x = u_r r_x + u_\theta \theta_x$$

$$u_y = u_r r_y + u_\theta \theta_y$$

$$u_{xx} = u_r r_{xx} + r_x (u_{rr} r_x + u_{r\theta} \theta_x) + u_\theta \theta_{xx} + \theta_x (u_{\theta r} r_x + u_{\theta\theta} \theta_x)$$

$$u_{yy} = u_r r_{yy} + r_y (u_{rr} r_y + u_{r\theta} \theta_y) + u_\theta \theta_{yy} + \theta_y (u_{\theta r} r_y + u_{\theta\theta} \theta_y)$$

Now, since  $r = (x^2 + y^2)^{1/2}$  and by part (b),

$$r_x = \frac{x}{(x^2 + y^2)^{1/2}} = \frac{x}{r} \quad r_{xx} = \frac{(x^2 + y^2)^{1/2}(1) - x \cdot \frac{x}{(x^2 + y^2)^{1/2}}}{x^2 + y^2} = \frac{r - \frac{x^2}{r}}{r^2} = \frac{r^2 - x^2}{r^3}$$

$$r_y = \frac{y}{(x^2 + y^2)^{1/2}} = \frac{y}{r} \quad r_{yy} = \frac{(x^2 + y^2)^{1/2}(1) - y \cdot \frac{y}{(x^2 + y^2)^{1/2}}}{(x^2 + y^2)^2} = \frac{r - \frac{y^2}{r}}{r^2} = \frac{r^2 - y^2}{r^3}$$

$$\theta_x = \frac{-y}{x^2 + y^2} = -\frac{y}{r^2} \quad \theta_{xx} = \frac{(x^2 + y^2)(0) - (-y)(2x)}{(x^2 + y^2)^2} = \frac{2xy}{r^4}$$

$$\theta_y = \frac{x}{x^2 + y^2} = \frac{x}{r^2} \quad \theta_{yy} = \frac{(x^2 + y^2)(0) - x(2y)}{(x^2 + y^2)^2} = -\frac{2xy}{r^4}$$

$$\Rightarrow u_{xx} + u_{yy} = u_r (r_{xx} + r_{yy}) + u_{rr} \left[ \frac{(r_x)^2 + (r_y)^2}{\frac{x^2}{r^2} + \frac{y^2}{r^2}} \right]$$

$$\frac{r^2 x^2 + r^2 y^2}{r^4} = \frac{2r^2 - r^2}{r^4} = \frac{1}{r}$$

$$+ 2u_{r\theta} (r_x \theta_x + r_y \theta_y) + u_{\theta\theta} \left[ \frac{(\theta_x)^2 + (\theta_y)^2}{\frac{y^2}{r^2} + \frac{x^2}{r^2}} \right]$$

$$-\frac{xy}{r^3} + \frac{xy}{r^3} = 0$$

$$\frac{2xy}{r^4} - \frac{2xy}{r^4} = 0$$

$$\frac{y^2}{r^4} + \frac{x^2}{r^4} = \frac{1}{r^2}$$

$$= \frac{1}{r} u_r + u_{rr} + \frac{1}{r^3} u_{\theta\theta}$$

$$= \frac{1}{r} \frac{\partial}{\partial r} (r u_r) + \frac{1}{r^3} u_{\theta\theta} \quad , \quad \text{since } \frac{1}{r} \frac{\partial}{\partial r} (r u_r) = \frac{1}{r} (r u_{rr} + u_r) = \frac{1}{r} u_r + u_{rr}$$

□

e) Use (d) to recheck that  $\ln r$  and  $\theta$  are harmonic functions.

Soln:  $\Delta(\ln r) = \frac{1}{r} \frac{\partial}{\partial r} \left[ r \frac{\partial}{\partial r} (\ln r) \right] + \frac{1}{r^3} \frac{\partial^2}{\partial \theta^2} (\ln r) = \frac{1}{r} \frac{\partial}{\partial r} \left( r \cdot \frac{1}{r} \right) + 0 = 0 = 0$

$$\Delta(\theta) = \frac{1}{r} \frac{\partial}{\partial r} \left[ r \frac{\partial}{\partial r} (\theta) \right] + \frac{1}{r^3} \frac{\partial^2}{\partial \theta^2} (\theta) = 0 + \frac{1}{r^3} \cdot 0 = 0 \quad \checkmark$$

2) Extended Maximum Principle in  $\mathbb{R}^2$

Let  $\Omega \subset \mathbb{R}^2$  be an open bounded domain with closure  $\bar{\Omega} = \Omega \cup \partial\Omega$

Let  $w \in C^2(\Omega)$  and continuous except at a finite number of boundary points  $P_1, \dots, P_k$ .

Let  $w$  be bounded,  $|w(x,y)| \leq m \quad \forall (x,y) \in \bar{\Omega}$

Let  $\Delta w \geq 0$  in  $\Omega$

Then if  $w \leq 0$  on  $\partial\Omega$ ,  $w \leq 0$  on  $\bar{\Omega}$ .

↑ we mean  $w \leq 0$  on  $\partial\Omega$ , except at  $P_1, P_2, \dots, P_k$  where any definition of  $w$  is irrelevant.

Prove this theorem.

Proof: We consider the function

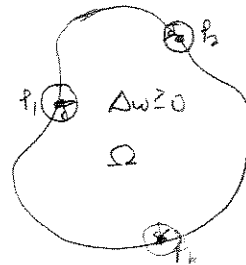
$$v = w + \varepsilon \sum_{j=1}^k \ln\left(\frac{\|x - P_j\|}{R}\right),$$

where  $\varepsilon > 0$  and  $R$  is such

that  $\Omega \subset \underset{\text{proper}}{B_R(P_j)} \quad \forall j$ , such an

$R$  exists because  $\Omega$  is bounded. We let  $v$  be defined on

a subdomain  $\Omega_\delta = \Omega \setminus \bigcup_{j=1}^k \bar{B}_\delta(P_j)$ , where  $\delta(\varepsilon) \rightarrow 0$ .



Now, for each  $j=1, \dots, k$  and  $\forall x \in \Omega$ ,  $\|x - P_j\| < R$ , so  $\ln\left(\frac{\|x - P_j\|}{R}\right) < 0$

$\forall x \in \Omega$  and  $\forall j$ .

Then  $v \leq w$  on  $\Omega_\delta$ . Additionally we have, by problem ① and the remarks that follow, that  $\ln\left(\frac{\|x - P_j\|}{R}\right) = \ln(\|x - P_j\|) - \ln(R)$  is

harmonic. Thus  $\Delta v = \Delta w \geq 0$  in  $\Omega_\delta$ .

Now, since  $w \leq 0$  on  $\partial\Omega \setminus \{P_1, \dots, P_k\}$ ,  $v \leq 0$  everywhere on

$\partial\Omega_\delta$  except possibly on  $\bigcup_{j=1}^k \partial B_\delta(P_j)$ . We will show that,

if  $\delta = R e^{-\frac{m}{\varepsilon}}$ , then in fact  $v \leq 0$  on all of  $\partial\Omega_\delta$ .

we have for  $x \in \bigcup_{j=1}^k \partial \bar{B}_\varepsilon(P_j)$

$$v = w + \varepsilon \sum_{j=1}^k \ln\left(\frac{\|x - P_j\|}{R}\right)$$

$$\leq M + \varepsilon \sum_{j=1}^k \ln\left(\frac{\|x - P_j\|}{R}\right)$$

$$\leq M + \varepsilon \ln\left(\frac{\|x - P_j\|}{R}\right) \quad \text{for some } 1 \leq j \leq k.$$

$$= M + \varepsilon \ln\left(\frac{\varepsilon}{R}\right)$$

$$= M + \varepsilon \ln\left(\frac{R e^{-\frac{M}{\varepsilon}}}{R}\right)$$

$$= M + \varepsilon \left(-\frac{M}{\varepsilon}\right)$$

$$= 0,$$

so  $v \leq 0$  on  $\partial \Omega_\varepsilon$ .

Thus, by the maximum principle,

$$v \leq 0 \text{ on } \bar{\Omega}_\varepsilon.$$

Fix  $x \in \bar{\Omega}_\varepsilon$ . Then, as  $\varepsilon \rightarrow 0$ ,  $\bar{\Omega}_\varepsilon \rightarrow \bar{\Omega}$ .

Also,  $\lim_{\varepsilon \rightarrow 0} v \leq \lim_{\varepsilon \rightarrow 0} 0 = 0$ . Since  $v = w + \varepsilon \sum_{j=1}^k \ln\left(\frac{\|x - P_j\|}{R}\right)$ ,

$v \xrightarrow[\text{as } \varepsilon \rightarrow 0]{} w$ . Thus  $w \leq 0$  on  $\bar{\Omega}$ .

③ a) Show that if

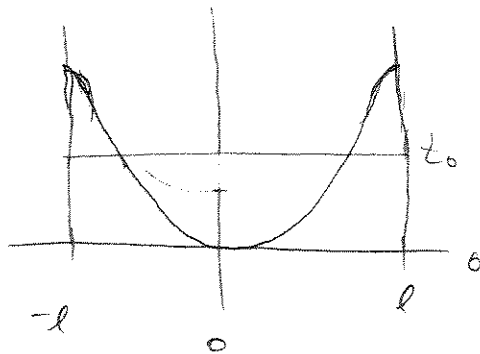
$$\frac{\partial u}{\partial t} - k \frac{\partial^2 u}{\partial x^2} = 0 \quad \text{for } 0 < x < l$$

$$\text{and } \frac{\partial u}{\partial x} \Big|_{x=0} = 0,$$

the maximum of  $u$  for  $0 \leq x \leq l$  and  $0 \leq t \leq \bar{t}$  must occur at

$$t=0 \text{ or } x=l$$

Soln: If we let  $u(x,t) = u(-x,t)$  for  $x \in [-l, 0]$ , then  $u(x,t)$  on  $-l \leq x \leq l$  is an even function such that  $u_t - k u_{xx} = 0$  for  $-l < x < l$ .



Let  $\bar{u}(x,t)$  be the extended function. Then  $\bar{u}(x,t)$  solves  $\bar{u}_t - k \bar{u}_{xx} = 0$   $-l < x < l$

$$\frac{\partial \bar{u}}{\partial x} \Big|_{x=0, t} = 0.$$

(This last condition holds because  $\bar{u}(x,t)$  is even about  $x=0$ .)  
By the maximum principle <sup>(\*) see below</sup>, the maximum value of  $\bar{u}$  must be attained on the parabolic boundary described by

$$(\partial \Omega_{\bar{u}} \times [0, t_0]) \cup (\Omega_{\bar{u}} \times \{0\})$$

where  $\Omega_{\bar{u}} = (-l, l) \times \{0\}$ .

Since  $u = \bar{u}|_{0 \leq x \leq l}$ , we conclude that the maximum of  $u$  must occur on  $([0, l] \times \{0\}) \cup (\{l\} \times [0, t_0])$ .

(\*) Indeed,  $\bar{u}$  must be bounded on  $-l \leq x \leq l \times 0 \leq t \leq t_0$  since  $\bar{u}$  is continuous and the set  $\{(x,t) \mid -l \leq x \leq l, 0 \leq t \leq t_0\}$  is compact. Clearly,  $\bar{u}$  is  $C^2$  on this set, since  $u$  is, i.e., since  $u$  is  $C^2$  at 0 from the right,  $\bar{u}$  will be  $C^2$  at 0 from the left, so  $\bar{u}$  will be  $C^2$  at 0. Thus, the maximum principle

(6)

3b) 13.4) Show that a soltn of the nonlinear eqn

$$\frac{\partial E}{\partial u}(\bar{x}, u(\bar{x}, t)) u_t - \operatorname{div}(K(\bar{x}, u(\bar{x}, t)) \nabla u) = 0$$

solves the max principle provided  $\frac{\partial E}{\partial u} > 0$  and  $K > 0$

proof:  $u(\bar{x}, t)$  is given.

$$\text{define } e(\bar{x}) := \frac{\partial E}{\partial u}(\bar{x}, u(\bar{x}, t))$$

$$\tilde{K}(\bar{x}) := K(\bar{x}, u(\bar{x}, t))$$

then  $u$  also solves the linear PDE

$$L(u) := \left[ e(\bar{x}) u_t - \nabla \tilde{K} \cdot \nabla u - K \Delta u = 0 \right] \quad e > 0.$$

consider

$$u^\varepsilon := u + \varepsilon \|x\|^2$$

$$\frac{\partial u^\varepsilon}{\partial t} = u_t; \quad \Delta u^\varepsilon = \Delta u + 2\varepsilon n.$$

consider the global maximum value of  $u^\varepsilon(\bar{x}, t)$  on the compact domain  $\bar{\Omega} \times [0, T]$ .

If it occurs on  $\Omega \times (0, T]$  then, say at  $(x_0, t_0)$ ,

$$\text{then } \nabla u^\varepsilon(x_0, t_0) = 0$$

$$u_t^\varepsilon(x_0, t_0) \geq 0 \quad (= 0 \text{ unless } t_0 = T)$$

$$u_{\xi\xi}^\varepsilon(x_0, t_0) \leq 0 \quad \text{for all spatial 2nd directional derivatives}$$

$$\Rightarrow \Delta u^\varepsilon(x_0, t_0) \leq 0$$

$$\Rightarrow L(u^\varepsilon)(x_0, t_0) = \underbrace{e(x) u_t^\varepsilon(x_0, t_0)}_{\geq 0} - \underbrace{\nabla \tilde{K} \cdot \nabla u^\varepsilon(x_0, t_0)}_{=0} - \underbrace{K \Delta u(x_0, t_0)}_{\geq 0}$$

$$\geq 0.$$

$$\text{but } L(u) = 0$$

$$\text{so } L(u^\varepsilon) = L(u) + \cancel{L(u)} + L(\varepsilon \|x\|^2) = 0 - K \cdot 2\varepsilon n < 0 \quad \neq 0.$$

Thus the max value of  $u^\varepsilon$  occurs only on the parabolic

$$\text{boundary } (\partial \Omega \times [0, T]) \cup (\Omega \times \{0\}) =: PB$$

So for fixed  $(\bar{x}, t) \in \bar{\Omega} \times [0, T]$

$$u(x, t) \leq u^\varepsilon(x, t) \leq \max_{PB} u^\varepsilon(x, t) \leq \max_{PB} u(x, t) + \varepsilon R^2$$

$$\lim_{\varepsilon \rightarrow 0} \Rightarrow u(x, t) \leq \max_{PB} u(x, t)$$

where  
 $\downarrow$   
 $\bar{\Omega} \subset B(0, R)$

4) If  $\vec{F} = \langle m(x,y), n(x,y) \rangle$  is an  $\mathbb{R}^2$  vector field, consider  $\mathbb{R}^2$  as the  $z=0$  plane in  $\mathbb{R}^3$ , extend  $\vec{F}$  as  $\langle m, n, 0 \rangle$  in  $\mathbb{R}^3$ . Let  $S = \Omega$  a domain in  $\mathbb{R}^2$  and check that Green's Theorem is a special case of Stokes' Theorem.

Soln: Stokes' Theorem states  $\iint_{\Omega} (\nabla \times \vec{F}) \cdot \hat{n} \, dA = \oint_{\partial \Omega} \vec{F} \cdot \hat{T} \, ds$ .

In this case  $\hat{n} = \hat{k}$ , and

$$\nabla \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \partial_x & \partial_y & \partial_z \\ m & n & 0 \end{vmatrix} = \hat{i}(-n_z) - \hat{j}(-m_z) + \hat{k}(n_x - m_y)$$

Since  $\vec{F} \in \mathbb{R}^2$ ,  $n_z = m_z = 0$ . Now,  $(\nabla \times \vec{F}) \cdot \hat{k} = n_x - m_y$ , so Stokes' Theorem becomes

$$\iint_{\Omega} (n_x - m_y) \, dA = \oint_{\partial \Omega} \vec{F} \cdot \hat{T} \, ds,$$

which is Green's Theorem.

b) We need to show Stokes' Theorem for the given graph by invoking Green's Theorem in  $D$ , i.e., we need to show that

$$\int_{\partial S} \vec{F} \cdot \hat{T} \, ds = \iint_S (\nabla \times \vec{F}) \cdot \hat{n} \, dS$$

We begin on the left. We can parameterize the curve  $\partial S$  by  $\vec{r}(t) = \langle x(t), y(t), f(x(t), y(t)) \rangle$ .

$$\text{Also, } \hat{T} = \frac{\vec{r}'(t)}{\|\vec{r}'(t)\|} = \frac{\langle x'(t), y'(t), f_x x'(t) + f_y y'(t) \rangle}{\|\vec{r}'(t)\|}$$

And, on  $\partial S$ ,  $\vec{F} = \langle m(x(t), y(t), f(x(t), y(t))), n, p \rangle$ ,

where  $N$  and  $P$  are also evaluated at  $\langle x(t), y(t), f(x(t), y(t)) \rangle$ .



Now,

$$\oint_{\partial D} \vec{F} \cdot \vec{T} \, ds = \int_t \langle M(x(t), y(t), f(x(t), y(t))), N, P \rangle \cdot \frac{\langle x', y', f_x x' + f_y y' \rangle}{\|\vec{r}'(t)\|} \|\vec{r}'(t)\| \, dt$$

since  $ds = \|\vec{r}'(t)\| \, dt$ .

The above equality is equal to

$$\int_t M(x(t), y(t), f(x(t), y(t))) x' + N y' + P(f_x x' + f_y y') \, dt.$$

Since  $x' dt = dx$ ,  $y' dt = dy$ , this becomes

$$\int_{\partial D} [M(x, y, f(x, y)) dx + N dy + P f_x dx + P f_y dy]$$

$$= \int_{\partial D} \left\{ [M(x, y, f(x, y)) + P f_x] dx + (N + P f_y) dy \right\}$$

By Green's  
theorem

$$\iint_D (N + P f_y)_x - [M(x, y, f(x, y)) + P f_x]_y \, dA$$

$$= \iint_D [N_x + N_z f_x + P f_{yx} + P_x f_y - [M_y + M_z f_y + P f_{xy} + P_y f_x]] \, dA$$

$$= \iint_D [-f_x (P_y - N_z) - f_y (M_z - P_x) + (N_x - M_y)] \, dA$$

$$= \iint_S \langle P_y - N_z, M_z - P_x, N_x - M_y \rangle \cdot \underbrace{\langle -f_x, -f_y, 1 \rangle}_{\hat{n}} \frac{1}{\sqrt{f_x^2 + f_y^2 + 1}} \, dS$$

$$\text{Now, } \nabla \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \partial_x & \partial_y & \partial_z \\ M & N & P \end{vmatrix} = \hat{i} (P_y - N_z) + \hat{j} (M_z - P_x) + \hat{k} (N_x - M_y),$$

so the above integral becomes

$$\iint_S (\nabla \times \vec{F}) \cdot \hat{n} \, dS. \quad \square$$