

Name \_\_\_\_\_

Math 3150 Problems  
Haberman Chapter H4

Due Date: Problems are collected on Wednesday.

## Chapter H4: 4.2, 4.3, 4.4 Vibrating String, Boundary Conditions, Fixed Ends BVP

### EXERCISES H4.2

#### Problem H4.2-1. (One-Dimensional String Derivation)

(a) Using the equation

$$\rho_0(x) \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2} + Q(x, t) \rho_0(x),$$

compute the sagged equilibrium position  $u_E(x)$  if  $Q(x, t) = -g$ . Use boundary conditions  $u(0) = 0$  and  $u(L) = 0$  with the equilibrium equation  $0 = T_0 \frac{d^2 u}{dx^2} - g \rho_0$  (formally,  $u_{tt} \equiv 0$ ).

(b) Show that  $v(x, t) = u(x, t) - u_E(x)$  satisfies

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}, \quad c^2 = \frac{T_0}{\rho_0(x)}.$$

Use superposition and the equation  $0 = u_{xx} - gc^2$ , valid for  $u = U_E$  according to part (a).

#### Problem H4.2-2. (Wave Speed $c$ )

Show that  $c^2 = \frac{T_0}{\rho_0}$  has the dimensions of velocity squared.

#### Problem XC-H4.2-3. (String Derivation)

Consider a particle whose x-coordinate (in horizontal equilibrium) is designated by  $\alpha$ . If its vertical and horizontal displacements are  $u$  and  $v$ , respectively, determine its position  $x$  and  $y$ . Then show that

$$\frac{dy}{dx} = \frac{\partial u / \partial \alpha}{1 + \partial v / \partial \alpha}.$$

#### Problem XC-H4.2-4. (String Derivation)

Derive equations for horizontal and vertical displacements without ignoring  $v$ . Assume that the string is perfectly flexible and that the tension is determined by an experimental law.

#### Problem XC-H4.2-5. (String Derivation, Constant Wave Speed)

Derive the partial differential equation for a vibrating string in the simplest possible manner. You may assume the string has constant mass density  $\rho_0$ , you may assume the tension  $T_0$  is constant, and you may assume small displacements (with small slopes).

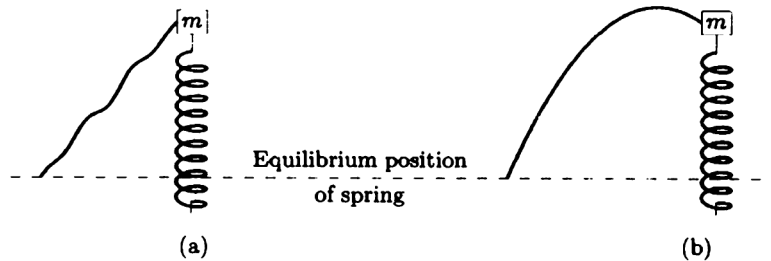
### EXERCISES H4.3

#### Problem XC-H4.3-1. (Boundary Conditions)

If  $m = 0$  in the model

$$m \frac{d^2 u}{dt^2}(0, t) = -k(u(0, t) - y_s(t) - L) + T_0 \frac{\partial u}{\partial x}(0, t) + g(t)$$

then which of the diagrams for the right end of the string shown in the figure can possibly be correct? *Briefly* explain. Assume that the mass can move only vertically.



### EXERCISES H4.4

#### Problem H4.4-1. (One-Dimensional String, Constant Wave Speed $c$ )

Consider vibrating strings of uniform density  $\rho_0$  and tension  $T_0$ . Submit only (a) and (b).

\* (a) What are the natural frequencies of a vibrating string of length  $L$  fixed at both ends?

\* (b) What are the natural frequencies of a vibrating string of length  $H$ , which is fixed at  $x = 0$  and "free" at the other end [i.e.,  $\partial u / \partial x(H, t) = 0$ ]? Sketch a few modes of vibration as in Fig. 1, H4.4.

(c) Show that the modes of vibration for the odd harmonics (i.e.,  $n = 1, 3, 5, \dots$ ) of part (a) are identical to modes of part (b) if  $H = L/2$ . Verify that their natural frequencies are the same. Briefly explain using symmetry arguments.

#### Problem H4.4-2. (Time-Dependent String Equation)

In Sec. H4.2 it was shown that the displacement  $u$  of a nonuniform string satisfies

$$\rho_0(x) \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2} + Q(x, t) \rho_0(x),$$

where  $Q$  represents the vertical component of the body force per unit length. If  $Q = 0$ , the partial differential equation is homogeneous. A slightly different homogeneous equation occurs if  $Q = \alpha u$ .

Submit only part (c).

(a) Show that if  $\alpha < 0$ , the body force is restoring (toward  $u = 0$ ). Show that if  $\alpha > 0$ , the body force tends to push the string further away from its unperturbed position  $u = 0$ .

(b) Separate variables if  $\rho_0(x)$  and  $\alpha(x)$  are assumed to depend on  $x$ , but  $T_0$  is constant for physical reasons. Analyze the time-dependent ordinary differential equation.

\* (c) Specialize part (b) to the constant coefficient case. Solve the initial value problem if  $\alpha < 0$ :

$$u(0, t) = 0, \quad u(L, t) = 0,$$

$$u(x, 0) = 0, \quad \frac{\partial u}{\partial t}(x, 0) = f(x).$$

What are the frequencies of vibration?

### Problem XC-H4.4-3. (Damped Vibrations of a One-Dimensional String)

Consider a slightly damped vibrating string that satisfies

$$\rho_0(x) \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2} - \beta \frac{\partial u}{\partial t}.$$

(a) Briefly explain why  $\beta > 0$ .

\* (b) Determine the solution (by separation of variables) that satisfies the boundary conditions

$$u(0, t) = 0 \text{ and } u(L, t) = 0$$

and the initial conditions

$$u(x, 0) = f(x) \text{ and } \frac{\partial u}{\partial t}(x, 0) = g(x).$$

You can assume that this frictional coefficient  $Q$  is relatively small ( $\beta^2 < 4\pi^2 \rho_0 T_0 / L^2$ ).

### Problem XC-H4.4-4. (Eigenfunction Expansion Method)

Redo Exercise 3(b), H4.4, by the eigenfunction expansion method.

### Problem XC-H4.4-5. (Damped Vibrations)

Redo Exercise 3(b), H4.4, if  $4\pi^2 \rho_0 T_0 / L^2 < \beta^2 < 16\pi^2 \rho_0 T_0 / L^2$ .

### Problem XC-H4.4-6. (d'Alembert Solution)

For the classical string vibration problem with clamped ends, use the series solution for  $u(x, t)$  to show that

$$u(x, t) = R(x - ct) + S(x + ct),$$

where  $R$  and  $S$  are some functions.

### Problem H4.4-7. (d'Alembert Solution)

If a vibrating string satisfying the one-dimensional string equation with fixed ends is initially at rest,  $g(x) = 0$ , with shape  $f(x)$  given, then show that

$$u(x, t) = \frac{1}{2} [F(x - ct) + F(x + ct)],$$

where  $F(x)$  is the odd periodic extension of  $f(x)$ .

Hints:

1. For all  $x$ ,  $F(x) = \sum_{n=1}^{\infty} A_n \sin(n\pi x/L)$ .

2.  $\sin a \cos b = \frac{1}{2} [\sin(a + b) + \sin(a - b)]$ .

Comment: This result shows that the practical difficulty of summing an infinite number of terms of a Fourier series may be avoided for the one-dimensional wave equation.

### Problem XC-H4.4-8. (d'Alembert Solution)

If a vibrating string satisfying the one-dimensional string equation with fixed ends is initially unperturbed,  $f(x) = 0$ , with the initial velocity  $g(x)$  given, then show that

$$u(x, t) = \frac{1}{2c} \int_{x-ct}^{x+ct} G(u) du,$$

where  $G(x)$  is the odd periodic extension of  $g(x)$ .

Hints:

1. For all  $x$ ,  $G(x) = \sum_{n=1}^{\infty} B_n \frac{n\pi c}{L} \sin(n\pi x/L)$ .

2.  $\sin a \sin b = \frac{1}{2} [\cos(a - b) - \cos(a + b)]$ .

See the comment after Exercise 7, H4.4.

### Problem XC-H4.4-9. (Energy Conservation)

From  $u_{tt} = c^2 u_{xx}$ , derive conservation of energy for a vibrating string,

$$\frac{dE}{dt} = c^2 u_x(x, t) u_t(x, t) \Big|_{x=0}^{x=L},$$

where the total energy  $E$  is the sum of the kinetic energy, defined by  $\int_0^L \frac{1}{2} (u_t)^2 dx$ , and the potential energy, defined by  $\int_0^L \frac{c^2}{2} (u_x)^2 dx$ .

**Problem XC-H4.4-10. (Total Energy)**

What happens to the total energy  $E$  of a vibrating string (see Exercise 9, H4.4)

- (a) If  $u(0, t) = 0$  and  $u(L, t) = 0$
- (b) If  $u_x(0, t) = 0$  and  $u(L, t) = 0$
- (c) If  $u(0, t) = 0$  and  $u_x(L, t) = -\gamma u(L, t)$  with  $\gamma > 0$
- (d) If  $\gamma < 0$  in part (c)

**Problem H4.4-11. (Potential and Kinetic Energies)**

Show that the potential and kinetic energies (defined in Exercise 9, H4.4) are equal for a traveling wave,  $u = R(x - ct)$ .

**Problem XC-H4.4-12. (Uniqueness)**

Using

$$\frac{dE}{dt} = c^2 u_x(x, t)u_t(x, t)\Big|_{x=0}^{x=L},$$

prove that the solution of the one-dimensional string equation with fixed ends is unique.

Remark. The result means that the solution can be computed with numerical software.

**Problem XC-H4.4-13. (Modes and Energies)**

(a) Using

$$\frac{dE}{dt} = c^2 u_x(x, t)u_t(x, t)\Big|_{x=0}^{x=L},$$

calculate the energy of one normal mode.

(b) Show that the total energy, when  $u(x, t)$  is a superposition of product solutions representing the solution, is the sum of the energies contained in each mode.