

Now $c_2 = 0$ since we do not want unbounded solutions for $r=0$.

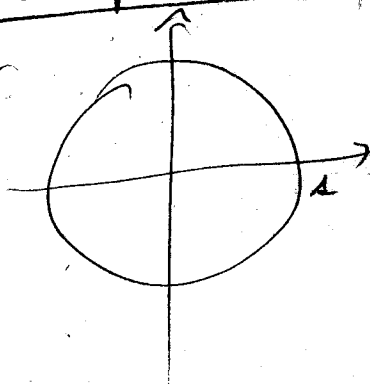
thus

$$u(r, \theta) = a_0 + \sum_{n=1}^{\infty} \left(\frac{r}{a}\right)^n [a_n \cos n\theta + b_n \sin n\theta]$$

$$\Rightarrow u(a, \theta) = f(\theta) = \sum_{n=1}^{\infty} a_n \cos n\theta + b_n \sin n\theta$$

$$\Rightarrow \left| \begin{array}{l} a_0 = \frac{1}{2\pi} \int_0^{2\pi} f(\theta) d\theta \\ a_n = \frac{1}{2\pi} \int_0^{2\pi} f(\theta) \cos n\theta d\theta \\ b_n = \frac{1}{2\pi} \int_0^{2\pi} f(\theta) \sin n\theta d\theta \end{array} \right.$$

Example 4.4.1



Steady state temp distrib on disk of radius 1 with $h(1, \theta) = f(\theta) = \begin{cases} 100 & 0 < \theta < \pi \\ 0 & \pi < \theta < 2\pi \end{cases}$

$$a_0 = \frac{1}{2\pi} \int_0^{\pi} 100 d\theta = 50$$

$$a_n = \frac{1}{\pi} \int_0^{\pi} 100 \cos n\theta d\theta = 0$$

$$b_n = \frac{1}{\pi} \int_0^{\pi} 100 \sin n\theta d\theta = -\frac{100}{n\pi} \cos n\theta \Big|_0^{\pi}$$

$$= \frac{100}{n\pi} (1 - \cos n\pi)$$

$$\Rightarrow u(r, \theta) = 50 + \frac{100}{\pi} \sum_{n=1}^{\infty} \frac{1}{n} (1 - \cos n\pi) r^n \sin n\theta$$

Of course by construction at $r=1$ we should get back Fourier series of f .

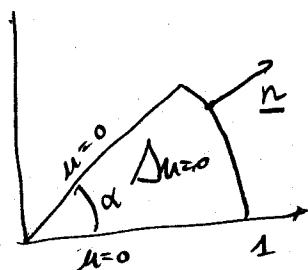
Other regions and Boundary conditions

(56)

With polar coordinates it is relatively easy to solve Laplace equation in annulus, wedges, or outside disk.

Example 4.4.4

Laplace equation in a wedge with Robin B.C.



$$\begin{cases} \Delta u = 0 & \text{in wedge} \\ u = 0 & \text{on sides of wedge} \\ \underline{n} \cdot \nabla u + u = -\theta & \text{on } r=1 \text{ boundary of wedge} \end{cases}$$

Recall Robin B.C. are in general:

$$\underline{n} \cdot \nabla u + a u = f$$

\underline{n} = unit outward pointing vector normal to boundary

In our case $\underline{n} \cdot \nabla u \Big|_{r=1} = \frac{\partial u}{\partial n} \Big|_{r=1} = \frac{\partial u}{\partial r}(1, \theta)$

Thus B.C. on part of wedge where $r=1$ is:

$$u_r(1, \theta) + u(1, \theta) + \theta = 0$$

Method: separation of variables

$$u(r, \theta) = R(r) \Theta(\theta)$$

as before we get (1) $\begin{cases} \Theta'' + \lambda^2 \Theta = 0 \\ \Theta(0) = \Theta(\alpha) = 0 \end{cases}$

and (2) $r^2 R'' + r R' - \lambda^2 R = 0$

Solving (1) we get $\Theta(\theta) = a \cos \lambda \theta + b \sin \lambda \theta$, but $\Theta(0) = 0 \Rightarrow a = 0$

and $\Theta(\alpha) = 0 \Rightarrow \lambda = \lambda_n = \frac{n\pi}{\alpha}$

$$\Rightarrow \Theta_n(\theta) = \sin \frac{n\pi}{\alpha} \theta$$

$n = 1, 2, \dots$

Now we get for (2):

$$r^2 R'' + r R' - \left(\frac{n\pi}{\alpha}\right)^2 R = 0$$

Solution is of the form $R(r) = a r^{\frac{n\pi}{\alpha}} + b r^{-\frac{n\pi}{\alpha}}$

But since $R(0)$ must be bounded we have $b = 0$

$$\Rightarrow R_n(r) = r^{\frac{n\pi}{\alpha}}$$

$$\text{and } u_n(r, \theta) = b_n r^{\frac{n\pi}{\alpha}} \sin\left(\frac{n\pi}{\alpha} \theta\right)$$

solves DE by construction.

\Rightarrow Solution is of the form:

$$u(r, \theta) = \sum_{n=1}^{\infty} b_n r^{\frac{n\pi}{\alpha}} \sin\left(\frac{n\pi}{\alpha} \theta\right), \quad \begin{matrix} r \in [0, 1], \\ \theta \in [0, \alpha]. \end{matrix}$$

To find b_n we look at B.C.:

$$u_r(0, \theta) = \sum_{n=1}^{\infty} b_n \frac{n\pi}{\alpha} r^{\frac{n\pi}{\alpha} - 1} \sin\left(\frac{n\pi}{\alpha} \theta\right)$$

$$u_r(1, \theta) = \sum_{n=1}^{\infty} b_n \frac{n\pi}{\alpha} \sin \frac{n\pi}{\alpha} \theta$$

$$\stackrel{\substack{\uparrow \\ \text{Robin B.C.}}}{=} -u(1, \theta) - \theta$$

$$= -\theta - \sum_{n=1}^{\infty} b_n \sin \frac{n\pi}{\alpha} \theta$$

$$\Rightarrow -\theta = \sum_{n=1}^{\infty} b_n \left(1 + \frac{n\pi}{\alpha}\right) \sin \frac{n\pi}{\alpha} \theta$$

Coeff in half range sine series of $-\theta$:

$$b_n \left(1 + \frac{n\pi}{\alpha}\right) = -\frac{2}{\alpha} \int_0^{\alpha} \theta \sin \frac{n\pi}{\alpha} \theta \cdot d\theta$$

$$= -\frac{2}{\alpha} \left[-\theta \frac{\alpha}{n\pi} \cos \frac{n\pi}{\alpha} \theta \Big|_0^{\alpha} + \frac{\alpha}{n\pi} \int_0^{\alpha} \cos \frac{n\pi}{\alpha} \theta \cdot d\theta \right]$$

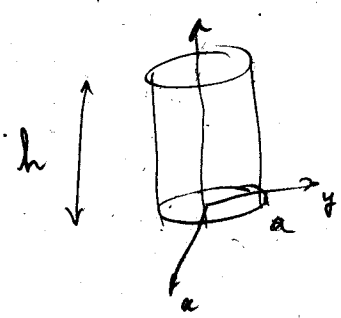
$$= \frac{2\alpha(-1)^n}{n\pi}$$

To conclude:

$$u(r, \theta) = \frac{2\alpha^2}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^n}{n(\pi+n\pi)} r^{\frac{n\pi}{\alpha}} \sin \frac{n\pi}{\alpha} \theta$$

§ 4.5 Laplace Equation in a Cylinder

Recall cylindrical coordinates (~ polar in xy plane + height z)



$$\begin{aligned} x &= r \cos \theta \\ y &= r \sin \theta \\ z &= z \end{aligned}$$

$$\begin{aligned} \Delta u &= u_{xx} + u_{yy} + u_{zz} \\ &= u_{rr} + \frac{1}{r} u_r + \frac{1}{r^2} u_{\theta\theta} + u_{zz} \end{aligned}$$

Laplace in polar coord.

The domain we work on is a cylinder defined by:
 $0 \leq \theta \leq 2\pi, 0 \leq r \leq a, 0 \leq z \leq h$.

Consider first the case where:

- bottom of cylinder is kept at temp 0
- lateral sides of cylinder
- top part of cylinder has a radially symmetric temperature

Then the PDE we would like to solve is:

$$\left\{ \begin{array}{ll} \Delta u = 0 & \text{inside cylinder} \\ u(r, 0) = 0 & 0 \leq r \leq a \\ u(a, z) = 0 & 0 \leq z \leq h \\ u(r, h) = f(r) & 0 \leq r \leq a \end{array} \right.$$