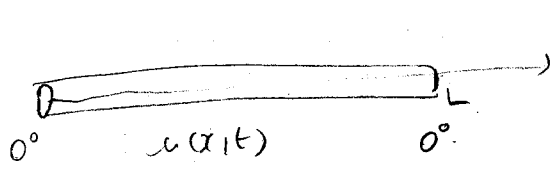


§ 3.5 The One-Dim. Heat eq (1DHEQ)



\$u(x,t)\$ = temperature distribution on a rod of length \$L\$, with ice baths on both ends.

\$u(x,t)\$ satisfies Heat Equation:

$$\begin{cases} u_t = c^2 u_{xx} \\ u(0,t) = u(L,t) = 0 & \text{B.C. = ice bath} \\ u(x,0) = f(x) & \text{I.C. = initial temperature distrib.} \end{cases}$$

Use method of separation of variable (it helps to see what we did last time in § 3.3)

Ansatz $u(x,t) = X(x)T(t)$

$$X T' = c^2 X'' T \Rightarrow \frac{T'}{c^2 T} = \frac{X''}{X} = k = \text{const indep of } x \text{ and } t.$$

$\underbrace{\hspace{1cm}}_{F(t)} \quad \underbrace{\hspace{1cm}}_{G(x)}$

We get 2 equations:

$$\begin{cases} X'' - kX = 0 \\ X(0) = X(L) = 0 \leftarrow \text{from B.C.} \end{cases}$$

and $T' - kc^2 T = 0$

$$k = -\mu^2, \mu_n = \frac{n\pi}{L}$$

$$T'_n + \left(\frac{cn\pi}{L}\right)^2 T_n = 0$$

$$X_n(x) = \sin\left(\frac{n\pi}{L} x\right), n=1,2,\dots$$

$$T_n(t) = b_n \exp\left[-\left(\frac{cn\pi}{L}\right)^2 t\right]$$

$$n=1,2,\dots$$

Thus we get a fundamental mode:

$$u_n(x,t) = b_n \sin\left[\frac{n\pi}{L} x\right] \exp\left[-\left(\frac{cn\pi}{L}\right)^2 t\right]$$

$$u(x,t) = \sum_{n=1}^{\infty} b_n \sin\left[\frac{n\pi}{L} x\right] \exp\left[-\left(\frac{cn\pi}{L}\right)^2 t\right]$$

1DHEQ is

homogeneous & lin.

Now what about initial conditions?

$$u(x,0) = f(x) = \sum_{n=1}^{\infty} b_n \sin \frac{n\pi}{L} x \quad \left| = \text{Sine Series of } f \right.$$

$$\Rightarrow b_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi}{L} x dx$$

exp. decay of temp

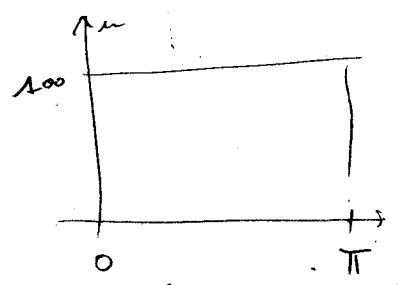
Summary

$$\begin{cases} u_t = c^2 u_{xx} \\ u(0,t) = u(L,t) = 0 \\ u(x,0) = f(x) \end{cases}$$

is solved by $u(x,t) = \sum_{n=1}^{\infty} b_n \sin \frac{n\pi}{L} x \exp\left[-\left(\frac{cn\pi}{L}\right)^2 t\right]$

where $b_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi}{L} x dx$

Example



$$\begin{cases} u_t = u_{xx} \\ u(0,t) = u(L,t) = 0 \\ u(x,0) = 100 \end{cases}$$

$$\Rightarrow u(x,t) = \sum_{n=1}^{\infty} b_n \sin nx \exp[-n^2 t]$$

$$u(x,t) = \sum_{k=0}^{\infty} \frac{e^{-(2k+1)^2 t}}{2k+1} \sin(2k+1)x$$

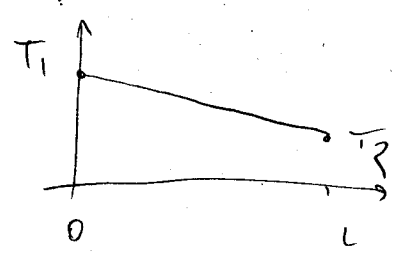
where $b_n = \frac{2}{\pi} \int_0^{\pi} 100 \sin nx dx$

$$= -\frac{200}{\pi n} \cos nx \Big|_0^{\pi} = \frac{200}{n\pi} (1 - (-1)^n)$$

Show matlab code

Steady state temp distrib:

$$u_t = 0 \Rightarrow u_{xx} = 0 \Rightarrow u(x) = Ax + B$$



$$\begin{aligned} u(x) &= T_1 \frac{L-x}{L} + T_2 \frac{x-0}{L} \\ &= \frac{T_2 - T_1}{L} x + T_1 \end{aligned}$$

Other Boundary Conditions :

$$(A) \begin{cases} u_t = c^2 u_{xx} \\ u(0,t) = T_1 \\ u(L,t) = T_2 \\ u(x,0) = f(x) \end{cases}$$

① find steady state

$$\Delta(x) = \frac{T_2 - T_1}{L} x + T_1$$

② shift solution by steady state

let $v = u - \Delta$ then

$$v_t = u_t$$

$$v_{xx} = u_{xx}$$

$\Rightarrow v$ solves:

$$(B) \begin{cases} v_t = c^2 v_{xx} \\ v(0,t) = 0 \\ v(L,t) = 0 \\ v(x,0) = f(x) - \Delta(x) \end{cases}$$

And we know how to solve (B):

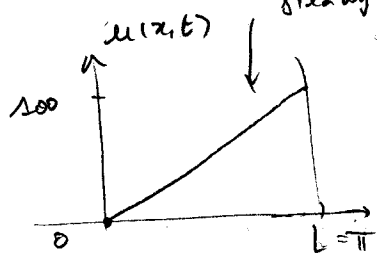
$$v(x,t) = \sum_{n=1}^{\infty} b_n e^{-\left(\frac{cn\pi}{L}\right)^2 t} \sin \frac{n\pi}{L} x$$

where $b_n = \frac{2}{L} \int_0^L (f(x) - \Delta(x)) \sin \frac{n\pi}{L} x \, dx$

③ go back to original problem

$$u(x,t) = v(x,t) + \Delta(x)$$

Example:



steady state temp $\Delta(x) = \frac{100}{\pi} x$

\Rightarrow I.C. for (B) is $100 - \frac{100}{\pi} x$

solve (B) for:

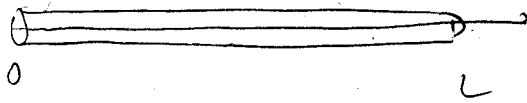
$$b_n = \frac{2}{\pi} \int_0^{\pi} \left(100 - \frac{100}{\pi} x\right) \sin nx \, dx$$

$$= \frac{200}{n\pi}$$

$$\Rightarrow u(x,t) = \Delta(x) + \sum_{n=1}^{\infty} b_n \sin nx e^{-n^2 t}$$

$$= \frac{100}{\pi} x + \frac{200}{\pi} \sum_{n=1}^{\infty} \frac{\sin nx}{n} e^{-n^2 t}$$

§ 3.6 Even more boundary conditions for 1D Heat Eq



metal rod with insulated ends
homogeneous Neumann B.C.

$$(1) \begin{cases} u_t = c^2 u_{xx}, & 0 < x < L, t > 0 \\ u_x(0, t) = u_x(L, t) = 0, & t > 0 \quad (\Leftrightarrow \text{no heat flux}) \\ u(x, 0) = f(x), & 0 < x < L \end{cases}$$

Solution using separation of variables.

$u(x, t) = X(x)T(t)$ = ansatz plug in in (1) gives:

$$\begin{cases} X'' - kX = 0 & T' - kc^2 T = 0 \\ X'(0) = X'(L) = 0 \end{cases}$$

$k = \mu^2 > 0$: $X(x) = a \cosh \mu x + b \sinh \mu x + B.C. \Rightarrow X(x) = 0$ (check)

$k = 0$: $X(x) = ax + b + B.C.$: $X'(x) = a$
 $X'(0) = X'(L) = 0 \Rightarrow a = 0$

\Rightarrow we get $X(x) = b$ a nontrivial solution

$k = -\mu^2 < 0$: $X(x) = a \cos \mu x + b \sin \mu x$

$X'(x) = -a \sin \mu x + b \cos \mu x$

$X'(0) = 0 \Rightarrow b = 0$

$X'(L) = 0 \Rightarrow a \sin \mu L = 0$

$\Rightarrow \mu = \mu_n = \frac{n\pi}{L}, n = 1, 2, \dots$

$\leadsto X_n(x) = \cos \frac{n\pi}{L} x, n = 1, 2, \dots$

Thus we get =

$T_0(t) = a_0$

$T_n(t) = a_n \exp \left[- \left(\frac{n\pi c}{L} \right)^2 t \right]$

Using superposition principle:

$$u(x,t) = a_0 + \sum_{n=1}^{\infty} a_n \exp\left[-\left(\frac{n\pi c}{L}\right)^2 t\right] \cos\left(\frac{n\pi}{L} x\right)$$

solves by construction (1)

What about initial conditions?

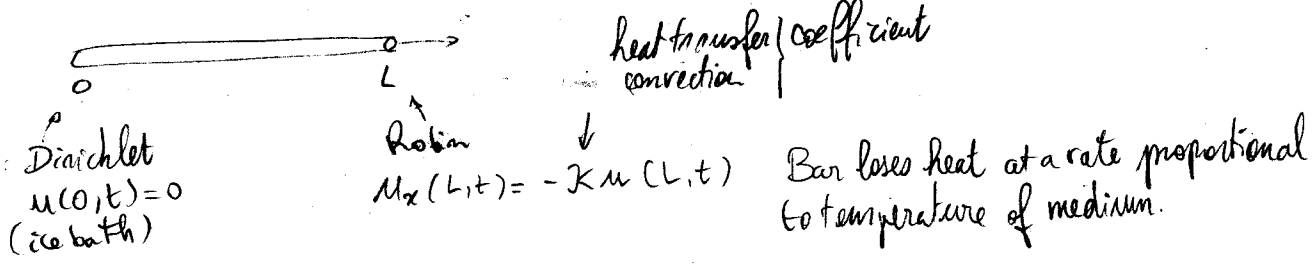
$$u(x,0) = f(x) = a_0 + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi}{L} x\right)$$

surprise! $a_n, n \geq 0$ are the coeff in the cosine series of $f(x)$.

$$\Rightarrow a_0 = \frac{1}{L} \int_0^L f(x) dx$$

$$a_n = \frac{2}{L} \int_0^L f(x) \cos\frac{n\pi}{L} x dx$$

Now more complicated Robin type boundary cdt:



$$(2) \begin{cases} u_t = c^2 u_{xx} & , 0 < x < L, t > 0 \\ u_x(0,t) = 0 & , t > 0 \\ u_x(L,t) = -K u(L,t) & , t > 0 \\ u(x,0) = f(x) & , 0 < x < L \end{cases}$$

Idea Use separation of variables to get:

$$\begin{cases} X'' - kX = 0 \\ X(0) = 0 & \text{(BC1)} \\ X'(L) = -K X(L) & \text{(BC2)} \end{cases}$$

$$T' - k^2 T = 0$$

$k = \mu^2 > 0$: $X(x) = a \cosh \mu x + b \sinh \mu x$. (BC1) $\Rightarrow a = 0$
 (BC2) $\Rightarrow b \frac{\sinh \mu L}{\cosh \mu L} = 0 \Rightarrow b = 0$
 $\Rightarrow X(x) = 0$ (trivial sol)

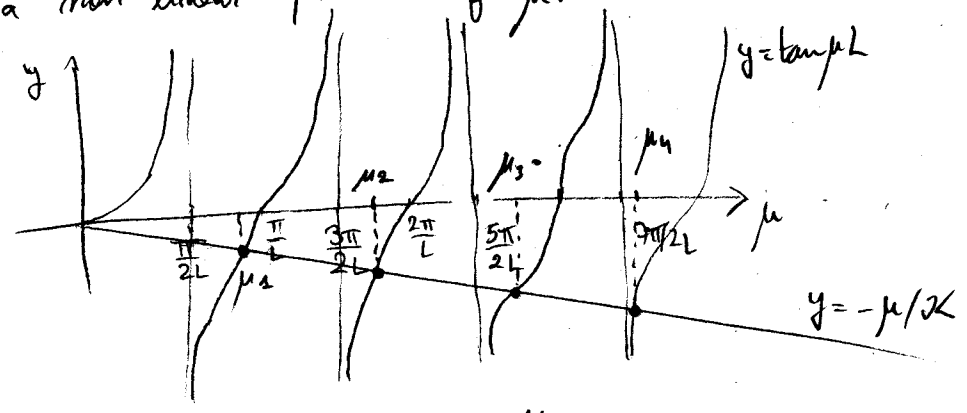
$k = 0$: $X(x) = ax + b$. (BC1) $\Rightarrow b = 0$
 (BC2) $\Rightarrow a = -K a L \Rightarrow a(1 + KL) = 0$
 $\Rightarrow a = 0$. (trivial sol)

$k = -\mu^2: X(x) = a \cos \mu x + b \sin \mu x$

(BC1) $\Rightarrow a = 0$

(BC2) $\Rightarrow \mu \cos \mu L = -K \sin \mu L \Leftrightarrow \tan \mu L = -\frac{\mu}{K}$

gives a non linear eq to solve for μ :



can solve (numerically) for roots $\mu_n, n \geq 1$

\Rightarrow get solutions:

$X_n(x) = \sin \mu_n x$

$T_n(t) = c_n \exp[-c^2 \mu_n^2 t]$

$\Rightarrow u(x,t) = \sum_{n=1}^{\infty} c_n \exp[-c^2 \mu_n^2 t] \sin \mu_n x$ solves (2) by construction

what about I.C.?

$u(x,0) = f(x) = \sum_{n=1}^{\infty} c_n \sin \mu_n x$

\sim "Generalized Fourier series" (Some series but not quite)

It turns out (Sturm-Liouville theory, § 6) that:

$\{ \sin \mu_n x \}_{n=1}^{\infty}$ form an orthogonal system of functions w.r.t inner product: $(u,v) = \int_0^L u(x)v(x) dx$

thus

$(f, \sin \mu_n x) = \sum_{n=1}^{\infty} c_n (\sin \mu_n x, \sin \mu_n x)$
 $= c_n (\sin \mu_n x, \sin \mu_n x)$

$\Rightarrow c_n = \frac{(f, \sin \mu_n x)}{(\sin \mu_n x, \sin \mu_n x)}$